

Diffusion and Homogenization limit of Boltzmann-Poisson system

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Outline

- 1 Introduction
- 2 Diffusion approximation
- 3 Diffusion & Homogenization
 - Linear free space
 - Remarks about convergence
 - Poisson Coupling & inflow boundary data
- 4 **Extensions & questions**

Fluid description: Drift-Diffusion equation

- $\rho := \rho(t, x)$: charge density

$$x \in \omega \subset \mathbb{R}^d, \quad d \geq 1, \quad t \geq 0$$

$$\partial_t \rho + \nabla_x \cdot j = 0$$

$$j = -D \nabla_x \rho + \mu \rho \nabla_x \phi, \quad D = \frac{k_B T}{e} \mu$$

$$-\Delta_x \phi = \rho$$

Kinetic description: Boltzmann Equation

- $f(t, x, \mathbf{v})dxdv$:= number of particles in the unit volume $dxdv$ around (x, \mathbf{v}) and at time t .

$$\frac{df}{dt} = \partial_t f + \mathbf{v} \cdot \nabla_x f + \frac{e}{m_p} \nabla_x \phi \cdot \nabla_v f := S + Q(f)$$

- **Electron-phonon collisions**

$$Q(f)(\mathbf{v}) = \int_{\mathbb{R}^d} \sigma(\mathbf{v}, \mathbf{v}') [M(\mathbf{v}) f(\mathbf{v}') - M(\mathbf{v}') f(\mathbf{v})] d\mathbf{v}'$$

$$M(\mathbf{v}) = \left(\frac{m}{2\pi k_B T} \right)^{d/2} \exp \left(\frac{-m|\mathbf{v}|^2}{2k_B T} \right)$$

Diffusion approximation

$$\varepsilon > 0 \quad " \quad t \rightarrow \varepsilon^2 t \quad x \rightarrow \varepsilon x \quad "$$

$$f^\varepsilon : f^\varepsilon(t, x, v) \geq 0, \quad (x, v) \in \Omega = \omega \times \mathbb{R}^d$$

$$\partial_t f^\varepsilon + \frac{1}{\varepsilon} (v \cdot \nabla_x f^\varepsilon - \nabla_x \Phi \cdot \nabla_v f^\varepsilon) - \frac{Q(f^\varepsilon)}{\varepsilon^2} = 0$$

$$Q(f) = \int_{\mathbb{R}^d} \sigma(v, v') [M(v) f(v') - M(v') f(v)] dv'$$

$$M(v) = \frac{e^{-|v|^2/2}}{(2\pi)^{d/2}}, \quad 0 < \sigma_1 \leq \sigma(v, v') = \sigma(v', v) \leq \sigma_2$$

Proposition (Dynamics of collisions)

- $-Q$: non-negative and selfadjoint on $L^2(dv/M(v))$
- $\mathcal{N}(Q) = \mathbb{R} M(v)$
- *Solvability condition*

$$\mathcal{R}(Q) = \mathcal{N}(Q)^\perp := \left\{ g = g(v) / \langle g, M \rangle = \int_{\mathbb{R}^d} g(v) dv = 0 \right\}$$

- Q is invertible on $\mathcal{N}(Q)^\perp$.
- *H-Theorem*

$$-\int_{\mathbb{R}^d} Q(f) \log \left(\frac{f}{M} \right) dv \geq \frac{1}{2} \int_{\mathbb{R}^d} \left(\sqrt{f} - \sqrt{\rho M} \right)^2 dv$$

Review on the question

- Poupaud (90') :

$$\|f^\varepsilon - \rho M(v)\|_{L^\infty(0,T; L^1)} \leq C_T \varepsilon$$

For : $\phi \in W_{t,x}^{2,\infty}$ & **well prepared** initial & inflow data

$$f^\varepsilon(t=0) = \rho_I(x)M(v)$$

$$f^\varepsilon = \rho_b(t,x)M(v), \quad (x,v) \in \Gamma^-,$$

$$\Gamma^\pm = \{(x,v) \in \partial\Omega, \quad \pm(v \cdot n(x)) > 0\}$$

Sketch of the proof

- 1 Regular Hilbert expansion \implies control en mass thanks to contraction property.
- 2 Analysis of the Half-Space problem associated to Q to correct the incoming boundary layer of order ε^{-1}

Main tool : • Entropy dissipation

- Good decay of $Q^{-1}(vM)$ as $|v| \gg 1$

References : Bardos, Degond, Perthame, Schmeiser,...

Diffusion & Homogenization

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- Free space : $(x, v) \in \mathbb{R}^{2d}$ & given potential

$$\Phi^\varepsilon := \Phi \left(t, x, \frac{x}{\varepsilon} \right)$$

Φ is periodic with respect to $y \in Y = [0, 1]^d$

$$\partial_t f^\varepsilon + \frac{1}{\varepsilon} (v \cdot \nabla_x f^\varepsilon - \nabla_x \Phi^\varepsilon \cdot \nabla_v f^\varepsilon) - \frac{Q(f^\varepsilon)}{\varepsilon^2} = 0$$

$$f^\varepsilon(t=0) = f_0^\varepsilon$$

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References: Goudon, Mellet, Poupaud, Zhang, ...

- The homogenized fluid system is associated with an effective Φ_e :

$$\Phi_e(t, x) = -\log \left(\int_Y e^{-\Phi(t, x, y)} dy \right)$$

Define

$$M_\Phi(t, x, y, v) = M(v) \exp(\Phi_e(t, x) - \Phi(t, x, y))$$

$$M_{\Phi^\varepsilon}(t, x, v) = M_\Phi \left(t, x, \frac{x}{\varepsilon}, v \right)$$

Double-scale expansion

$$f^\varepsilon = f_0(t, x, x/\varepsilon, v) + \varepsilon f_1(t, x, x/\varepsilon, v) + \varepsilon^2 f_2(t, x, x/\varepsilon, v) + \dots$$

$f_i : Y = [0, 1]^d$ –periodic with respect to the fast variable.

$$y := x/\varepsilon \quad \nabla_x f_i \longrightarrow \nabla_x f_i + \frac{1}{\varepsilon} \nabla_y f_i$$

Identification

$$\varepsilon^{-2} : \quad Lf_0 := v \cdot \nabla_y f_0 - \nabla_y \Phi \cdot \nabla_v f_0 - Q(f_0) = 0$$

$$\varepsilon^{-1} : \quad Lf_1 = -v \cdot \nabla_x f_0 + \nabla_x \Phi \cdot \nabla_v f_0$$

$$\varepsilon^0 : \quad Lf_2 = -\partial_t f_0 - v \cdot \nabla_x f_1 + \nabla_x \Phi \cdot \nabla_v f_1$$

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 References: Allaire, E, Ngutseng,...

Proposition (Cell problem)

- L is maximal monotone on $L^2_{M_\Phi} := L^2(dydv/M_\Phi)$
- $\mathcal{N}(L) = \mathbb{R} M_\Phi$
- $\mathcal{R}(L) = \left\{ g = g(y, v) \in L^2_{M_\Phi} / \int \int_{Y \times \mathbb{R}^d} g \, dydv = 0 \right\}$

$$\forall g \in \mathcal{R}(L) \quad \exists f \in D(L) / Lf = g$$

$$f \text{ is unique iff } \int \int_{Y \times \mathbb{R}^d} f \, dydv = 0$$

$$f_0 = \rho(t, x) M(v) \exp(\Phi_e(t, x) - \Phi(t, x, y)) := \rho(t, x) M_\Phi$$

$$f_1 = -(\nabla_x \rho + \rho \nabla_x \Phi_e)(t, x) \cdot L^{-1}(v M_\Phi)(t, x, y)$$

Homogenized fluid system:

$$\partial_t \rho + \nabla_x \cdot \mathbf{J}(\rho, \Phi_e) = 0$$

$$\mathbf{J}(\rho, \Phi_e) = -\mathbb{D}(t, x) [\nabla_x \rho + \rho \nabla_x \Phi_e]$$

$$\mathbb{D}(t, x) = \int \int_{Y \times \mathbb{R}^d} v \otimes L^{-1}(v M_\Phi) dy dv$$

Remark

- Entropy inequality $\implies \mathbb{D}$ is nonnegative
- L is not selfadjoint $\implies \mathbb{D}$ is not necessary symmetric

Uniform estimates

- $x \in \mathbb{R}^d \implies$ No boundary effect
- Total conservation of mass :

$$\frac{d}{dt} \|f^\varepsilon\|_{L^1} = 0$$

- L^∞ – maximum principle: If

$$f^\varepsilon(t=0) \lesssim M(v)$$

\implies

$$f^\varepsilon \lesssim \exp\left(-\frac{|v|^2}{2} + \int_0^T \|\partial_t \Phi\|_{L^\infty}(s) ds\right)$$

Relative Entropy

$$(v \cdot \nabla_x - \nabla_x \Phi^\varepsilon \cdot \nabla_v) e^{-|v|^2/2 - \Phi^\varepsilon} = Q(e^{-|v|^2/2 - \Phi^\varepsilon}) = 0$$

Multiplying by $f^\varepsilon / e^{-|v|^2/2 - \Phi^\varepsilon}$, and integrate by parts :

$$\frac{d}{dt} \int_{\mathbb{R}^{2d}} \left[\frac{f^\varepsilon{}^2}{e^{-|v|^2/2 - \Phi^\varepsilon}} \right] + \frac{\sigma_1}{\varepsilon^2} \int_{\mathbb{R}^{2d}} \frac{(f^\varepsilon - \rho^\varepsilon M)^2}{e^{-|v|^2/2 - \Phi^\varepsilon}} \leq 0$$



$$\rho^\varepsilon \quad \& \quad j^\varepsilon := \frac{1}{\varepsilon} \int_{\mathbb{R}^d} f^\varepsilon v \, dv \quad \in_b \quad L_{loc}^2(\mathbb{R}^+, L_x^2)$$

⇒ We can deal with double scale convergence

Denote :

- $\bar{f}, \bar{\rho}, \bar{j} \dots$: double-scale limits of $f^\varepsilon, \rho^\varepsilon, j^\varepsilon, \dots$

$$\varepsilon^2 \partial_t f^\varepsilon + \varepsilon v \cdot \nabla_x f^\varepsilon - \varepsilon \nabla_x \Phi^\varepsilon \cdot \nabla_v f^\varepsilon - Q(f^\varepsilon) = 0$$

$$\implies f^\varepsilon \xrightarrow{2-s} \bar{f} := \rho(t, x) M_{\Phi_H} \in \mathcal{N}(L)$$

- Entropy-Dissipation \implies

$$f^\varepsilon = \rho^\varepsilon M(v) + O(\varepsilon)_{L^2}$$

- Expression of $j(t, x)$?

$$g^\varepsilon = \frac{1}{\varepsilon} (f^\varepsilon - \rho^\varepsilon M) \in_b L^2 \xrightarrow{2-s} \bar{g} \in L^2_{t,x,y,v}$$

Green formula with $\frac{\varepsilon \psi(t, x, x/\varepsilon, v)}{M_{\Phi^\varepsilon}} / \nabla_y \cdot \int_{\mathbb{R}^d} v \psi dv = 0$

$$j^\varepsilon \rightharpoonup J(\rho, \Phi_e)$$

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- **Remark 1** : Compacity method ?
 - 1 NO strong convergence in L^p & velocity averaging lemma
 - 2 div-curl Lemma FAILS
- **Remark 2** : L^1 -contraction ?

$$f^\varepsilon(t=0) = \rho(t, x) M_{\Phi^\varepsilon}(t=0) : \quad \text{well prepared}$$

$$r^\varepsilon = f^\varepsilon - (f_0 + \varepsilon f_1 + \varepsilon^2 f_2)(t, x, x/\varepsilon, v)$$

$$\begin{cases} \partial_t r^\varepsilon + \frac{1}{\varepsilon} (v \cdot \nabla_x r^\varepsilon - \nabla_x \Phi^\varepsilon \cdot \nabla_v r^\varepsilon) - \frac{Q(r^\varepsilon)}{\varepsilon^2} = -\varepsilon \partial_t f_1 - \varepsilon^2 \partial_t f_2 + \dots \\ r^\varepsilon(t=0) = -\varepsilon f_1 - \varepsilon^2 f_2 \end{cases}$$

NEED : Good decay on the derivatives of $\chi = L^{-1}(vM)$

$$v \cdot \nabla_y \chi - \nabla_y \Phi \cdot \nabla_v \chi - Q(\chi) = v M_\Phi$$

\Rightarrow : Difficulty even for linear free case !!

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Poisson Coupling & inflow boundary data ??

(ε : fixed)

- If we regularize the Boltzmann-Poisson system we can only establish uniform estimate :

$$f \in L \log L, \quad \nabla_x \Phi \in L^2$$

This DOES NOT give a sense for $\nabla_x \Phi \cdot \nabla_v f$ in \mathcal{D}' !

- L^∞ estimate on f need $\partial_t \Phi \in L^\infty$!!
- Entropy-dissipation gives only $j \in L^1$ which does not imply $\partial_t \Phi \in L^\infty$... "NO L^p estimate" !!
- $j \in L^1 \not\Rightarrow \partial_t \nabla_x \Phi \in L^1_{loc}$!

$(\varepsilon \rightarrow 0)$

- Entropy production terms due to the boundary and no direct conservation of mass
- We can not use the Hilbert expansion since the Boltzmann equation has no 'weak PDE sense'.
- Assume that we can use the Hilbert expansion...we need to control $\varepsilon \partial_t \nabla_x \Phi^\varepsilon$ in L^1 !! but for fixed ε

$$-\Delta_x \partial_t \Phi^\varepsilon = \partial_t \rho^\varepsilon = \nabla \cdot j^\varepsilon \in W_x^{1,1} \quad \not\Rightarrow \partial_t \nabla_x \Phi^\varepsilon \in L_{loc}^1$$

- We need strong convergence on f^ε and the self consistent potential ! we need to pass to the limit in some $\beta(f^\varepsilon)$...

How to deal ??

For only macroscopic variation

- with N. Ben Abdallah : **1D case**

We constructed a weak solution of the 1D Boltzmann-Poisson system. We used a **Hybrid Hilbert expansion** and the relative entropy to analyze the entropy production terms due to the boundary and obtain in a first step uniform estimates. Then we pass to the limit like Poupaud and exhibit a rate convergence in ε .

- with N. Masmoudi : **multi-D case**

We prove existence of renormalized solution and prove strong convergence using a "limiting" L^1 -velocity averaging but without rate of convergence

$$f^\varepsilon := f^\varepsilon(t, x, v), \quad (x, v) \in \Omega = \omega \times \mathbb{R}^d$$

$$\left\{ \begin{array}{l} \partial_t f^\varepsilon + \frac{1}{\varepsilon} (v \cdot \nabla_x f^\varepsilon - \nabla_x \Phi_T^\varepsilon \cdot \nabla_v f^\varepsilon) - \frac{Q(f^\varepsilon)}{\varepsilon^2} = 0 \\ \Phi_T^\varepsilon(t, x) = \Phi_P^\varepsilon(t, x) + \Phi_H(x, \frac{x}{\varepsilon}) \\ f^\varepsilon|_{\Gamma^-} = \rho b(t, x) M_{\Phi_H}(x, \frac{x}{\varepsilon}) \quad (\Gamma^- : v \cdot n(x) < 0), \\ f^\varepsilon|_{t=0} = f_I^\varepsilon \end{array} \right.$$

$$-\Delta_x \Phi_P^\varepsilon = \rho^\varepsilon = \int_{\mathbb{R}^d} f^\varepsilon dv, \quad (\Phi_P^\varepsilon)|_{\partial\omega} = 0$$

$$\int_{\Omega} f_l^\varepsilon (1 + |v|^2 + |\log f_l^\varepsilon|) \leq C \quad \& \quad \Phi_H \in W^{2,\infty}(\bar{\omega} \times Y)$$

Theorem

$$f^\varepsilon e^{\Phi_H(x, x/\varepsilon)} \rightarrow \rho M(v) e^{\Phi_e(x)} \quad \text{in } L_t^1(L_{x,v}^1)$$

$$\nabla_x \Phi_P^\varepsilon \rightarrow \nabla_x \Phi_P \quad \text{in } L_t^2(L_x^p), \quad p < 2$$

(ρ, Φ_P) : solution of the homogenized Drift-Diffusion-Poisson system associated with $\Phi_P + \Phi_e$.

$$\tilde{f}^\varepsilon = e^{\Phi_H(x, x/\varepsilon)} f^\varepsilon$$

Definition & Theorem ε : fixed $\exists (f^\varepsilon, \Phi_P^\varepsilon) /$

1. $\forall \beta \in C^1(\mathbb{R}^+), |\beta(t)| \leq C(\sqrt{t} + 1) \beta' \in L^\infty$

$$\partial_t \beta(\tilde{f}^\varepsilon) + \frac{1}{\varepsilon} \left[v \cdot \nabla_x \beta(\tilde{f}^\varepsilon) - \nabla_v \cdot (\nabla_x \Phi_T^\varepsilon \beta(\tilde{f}^\varepsilon)) \right] = \beta'(\tilde{f}^\varepsilon) \frac{Q(\tilde{f}^\varepsilon)}{\varepsilon^2}$$

2. $\forall \lambda > 0, \theta_{\varepsilon, \lambda} = \sqrt{f^\varepsilon + \lambda M e^{-\Phi_H(x, x/\varepsilon)}}$

$$\begin{aligned} \varepsilon \partial_t \theta_{\varepsilon, \lambda} + v \cdot \nabla_x \theta_{\varepsilon, \lambda} - \nabla_v \cdot (\nabla_x \Phi_T^\varepsilon \theta_{\varepsilon, \lambda}) - \frac{Q(\tilde{f}^\varepsilon)}{2 \varepsilon \theta_{\varepsilon, \lambda}} \\ = \frac{\lambda}{2 \theta_{\varepsilon, \lambda}} v M \cdot \nabla_x \Phi_P^\varepsilon e^{-\Phi_H(x, x/\varepsilon)} \end{aligned}$$

3. Continuity equation

$$\partial_t \rho^\varepsilon + \nabla_x \cdot j^\varepsilon = 0$$

4. Entropy inequality

$$\left[\int_{\Omega} f^\varepsilon \left(\Phi_H^\varepsilon + \frac{|v|^2}{2} + \log f^\varepsilon \right) + \frac{1}{2} \|\nabla_x \Phi_P^\varepsilon\|_{L^2}^2 \right]_0^t - \frac{1}{\varepsilon^2} \int_0^t \int_{\Omega} Q(f^\varepsilon) \log \frac{f^\varepsilon}{M}$$

$$\leq -\frac{1}{\varepsilon} \int_0^t \int_{\Gamma^+ \cup \Gamma^-} f^\varepsilon \left(2\Phi_H^\varepsilon + \frac{|v|^2}{2} + \log f^\varepsilon \right) (v \cdot n(x))$$

Proposition

- $$\mathcal{M}^\varepsilon(t) + \mathcal{K}^\varepsilon(t) + \|\nabla_x \Phi_P^\varepsilon(t)\|_{L^2} + \frac{\mathcal{R}^\varepsilon(t)}{\varepsilon^2} + \int_0^t \|j^\varepsilon(s)\|_{L^1}^2 ds$$

$$\in_b L_{loc}^\infty(\mathbb{R}^+)$$

$$\mathcal{M}^\varepsilon(t) = \|f^\varepsilon(t)\|_{L^1(\Omega)} \quad \mathcal{K}^\varepsilon(t) = \left\| \frac{|v|^2}{2} f^\varepsilon(t) \right\|_{L^1(\Omega)}$$

$$\mathcal{R}^\varepsilon(t) = \int_0^t \int_\Omega \left(\sqrt{f^\varepsilon} - \sqrt{\rho^\varepsilon M} \right)^2$$

$$2\alpha\mathcal{M}^\varepsilon(t) + \mathcal{K}^\varepsilon(t) + \|\nabla_x \Phi_P^\varepsilon(t)\|_{L^2}^2 + \frac{\mathcal{R}^\varepsilon(t)}{\varepsilon^2} \\
\leq C_T - \frac{1}{\varepsilon} \int_0^t \int_{\Gamma^+} |\mathbf{v} \cdot \mathbf{n}(x)| [f^\varepsilon(\mathbf{v}) - f^\varepsilon(-\mathbf{v})] (1 + \mathcal{E}_F)$$

$\mathcal{E}_F(t, x) = \log \left(\frac{\rho_b(t, x)}{(2\pi)^{d/2}} \right)$: macroscopic quasi-Fermi level

$$(1 + \mathcal{E}_F(t, x))(\partial_t \rho^\varepsilon + \nabla_x \cdot \mathbf{j}^\varepsilon) = 0$$

$$\left| \frac{1}{\varepsilon} \int_0^t \int_{\Gamma^+} |\mathbf{v} \cdot \mathbf{n}(x)| [f^\varepsilon(\mathbf{v}) - f^\varepsilon(-\mathbf{v})] (1 + \mathcal{E}_F) \right| \leq C_T + \alpha \mathcal{M}^\varepsilon(t) + \frac{\mathcal{R}^\varepsilon(t)}{2\varepsilon^2} \\
+ \int_0^t \{\mathcal{M}^\varepsilon + \mathcal{K}^\varepsilon\}(s) ds$$

⇒

GRONWALL.

- $f^\varepsilon \in_b L \text{Log} L$, $\nabla_x \Phi_P^\varepsilon \in_b L_x^2$
- f^ε , $f_{|\Gamma^+}^\varepsilon$ and ρ^ε w. r. compact in L^1
- $\frac{Q(f^\varepsilon)}{\varepsilon}$ w. r. compact in L^1 (only $L^1!$)
- $\left. \begin{array}{l} \nabla \Phi_P^\varepsilon \in L^2, \quad \Delta \Phi_P^\varepsilon \in L^1 \\ \partial_t \nabla \Phi_P^\varepsilon := \nabla \Delta^{-1} \nabla \cdot j^\varepsilon \in L_t^2(W_{loc}^{-s,p}) \\ p > 1, s > d - d/p \end{array} \right\} \implies \left. \begin{array}{l} \nabla_x \Phi_P^\varepsilon \text{ cpt in } L_t^2(L_x^p) \\ \forall p < 2. \end{array} \right.$

Velocity averaging lemma

Lemma

$h^\varepsilon \in_b L^2$, $h_0^\varepsilon \in_b L^1$ and $h_1^\varepsilon \in [L^1_{t,x}(L^1_{loc,v})]^d$ bounded

$$\varepsilon \partial_t h^\varepsilon + v \cdot \nabla_x h^\varepsilon = h_0^\varepsilon + \nabla_v \cdot h_1^\varepsilon$$

$$\rho_\psi^\varepsilon := \int_{\mathbb{R}^d} h^\varepsilon(t, x, v) \psi(v) dv \quad \text{for } \psi \in \mathcal{D}(\mathbb{R}_v^d)$$

Then,

$$\lim_{z \rightarrow 0} \left(\sup_{\varepsilon < 1} \|\rho_\psi^\varepsilon(t, x + z) - \rho_\psi^\varepsilon(t, x)\|_{L^1_{t,x}} \right) = 0$$

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 References: Golse, Jabin, Lions, St. Raymond, Tadmor,...

Application:

$$\beta(\mathbf{s}) = \frac{\mathbf{s}}{1+\mathbf{s}}, \quad \beta_\delta(\mathbf{s}) = \frac{\mathbf{s}}{1+\delta\mathbf{s}}$$

- 1 $0 \leq \beta_\delta \leq \min(1, 1/\delta)$
- 2 $|\beta_\delta(\mathbf{s}) - \beta_\delta(\mathbf{t})| \leq \min(|\mathbf{s} - \mathbf{t}|, |\sqrt{\mathbf{s}} - \sqrt{\mathbf{t}}|/\delta)$
- 3 $|\beta'_\delta(\mathbf{s}) - 1| \leq 2\delta\mathbf{s}$
- 4 $|\mathbf{s}\beta'_\delta(\mathbf{s}) - \mathbf{t}\beta'_\delta(\mathbf{t})| \leq \min\left[(1 + \delta)|\mathbf{t} - \mathbf{s}|, C_\delta|\sqrt{\mathbf{s}} - \sqrt{\mathbf{t}}|\right]$

δ fixed $h^\varepsilon := \beta_\delta(\tilde{f}^\varepsilon) \in_b L^2$

satisfies the lemma with

$$h_0^\varepsilon : = \left[\frac{Q(\tilde{f}^\varepsilon)}{\varepsilon} + v \cdot (\nabla_x \Phi_H)_\varepsilon \tilde{f}^\varepsilon \right] \beta'_\delta(f^\varepsilon) \\
 + \frac{v \cdot (\nabla_y \Phi_H)_\varepsilon}{\varepsilon} \left[\tilde{f}^\varepsilon \beta'_\delta(\tilde{f}^\varepsilon) - \tilde{\rho}^\varepsilon M \beta'(\tilde{\rho}^\varepsilon M) \right] \in_b L^1$$

$$h_1^\varepsilon = \left[\nabla_x \Phi_P^\varepsilon + \nabla_x \Phi_H^\varepsilon \right] \beta_\delta(\tilde{f}^\varepsilon) + (\nabla_x \Phi_H)_\varepsilon \left(\frac{\beta_\delta(\tilde{f}^\varepsilon) - \beta_\delta(\tilde{\rho}^\varepsilon M)}{\varepsilon} \right) \\
 \in_b L^1_{t,x}(L^2_{loc,v})$$

Lemma $\implies \int_{\mathbb{R}^d} \beta_\delta(\tilde{f}^\varepsilon) \psi dv$ bounded in $L^1(0, T; cpt(L^1_x))$

$$\sup_{\varepsilon < 1} \|\beta_\delta(\tilde{f}^\varepsilon) - \tilde{f}^\varepsilon\|_{L^1(0, T; L^1)} \rightarrow 0 \quad \text{as } \delta \rightarrow 0$$

- $\delta \rightarrow 0 \Rightarrow \int_{\mathbb{R}^d} \tilde{f}^\varepsilon \psi \, dv$ bounded in $L^1(0, T; \text{cpt}(L^1_x))$
- Kinetic Energy \Rightarrow we can take $\psi = 1$

$$\tilde{\rho}^\varepsilon : \text{bounded in } L^1(0, T; \text{cpt}(L^1_x))$$

We have x -compactness for the modified density $\tilde{\rho}^\varepsilon$:

- $\partial_t \rho^\varepsilon = -\nabla_x \cdot j^\varepsilon$: bounded in $L^2(0, T; W_x^{-1,1})$
 $\implies \tilde{\rho}^\varepsilon$ compact in $L^1((0, T) \times \omega)$.

- **Remark 1.** We only use time-compactness of ρ^ε because

$$\partial_t \tilde{\rho}^\varepsilon + \nabla_x \cdot \tilde{j}^\varepsilon = \left(\nabla_x \Phi_H + \frac{1}{\varepsilon} \nabla_y \Phi_H \right)_\varepsilon \cdot \tilde{j}^\varepsilon$$

It is not clear how we can obtain directly the time-compactness for $\tilde{\rho}^\varepsilon$.

- **Remark 2.** We can not use the Div-curl lemma since $\tilde{\rho}^\varepsilon$ and \tilde{j}^ε are only L^1

Consequences:

$$\exists \rho \in L^1_{loc}(\mathbb{R}^+; L^1_x) \text{ \& } \Phi_P \in L^\infty_{loc}(\mathbb{R}^+; H^1_0(\omega))$$

$$\textcircled{1} \quad \|\rho^\varepsilon - \rho e^{\Phi_\varepsilon(x) - \Phi_H(x, x/\varepsilon)}\|_{L^1_{t,x}} \rightarrow 0$$

$$\textcircled{2} \quad \rho^\varepsilon \text{ converges weakly in } L^1_{t,x} \text{ to } \rho$$

$$\textcircled{3} \quad \sqrt{f^\varepsilon} \text{ \& } \sqrt{\rho^\varepsilon M} \xrightarrow{2-s \text{ strongly}} \sqrt{\rho M_{\Phi_H}} \text{ in } L^2_{t,x,v}$$

$$\textcircled{4} \quad -\Delta_x \Phi_P = \rho,$$

$$\nabla_x \Phi_P^\varepsilon \rightarrow \nabla_x \Phi_P \text{ in } L^2_{loc}(\mathbb{R}^+; L^p_x), \quad \forall p < 2$$

USE equation on $\sqrt{f^\varepsilon + \lambda M e^{-\Phi_H(x, x/\varepsilon)}}$ with an oscillating test function $\psi(t, x, x/\varepsilon, v)$ to identify the weak limit of j^ε

(ρ, ϕ) satisfies the homogenized system

$$\partial_t \rho + \nabla_x \cdot \mathbf{J}(\rho, \Phi_P + \Phi_e) = 0,$$

$$\mathbf{J}(\rho, \phi) = \sqrt{\rho} [-\mathbb{D}(x)(2\nabla_x \sqrt{\rho} + \nabla_x(\Phi_P + \Phi_e) \sqrt{\rho})],$$

$$\mathbb{D} = - \int_{Y \times \mathbb{R}^d} (\mathbf{v} \otimes L^{-1}(\mathbf{v} M_{\Phi_H})) dy dv > 0$$

$$L = \mathbf{v} \cdot \nabla_y - \nabla \Phi_H \cdot \nabla_v - Q$$

$$M_{\Phi_H} = M(\mathbf{v}) e^{\Phi_e(x) - \Phi_H(x,y)}$$

Moreover,

$$-\Delta \phi = \rho,$$

$$\rho|_{t=0} = \lim \int_{\mathbb{R}^d} f_0^\varepsilon dv \quad (L^1 - w),$$

$$\rho|_{\partial\omega} = \rho b$$

Regularity of $J(\rho, \Phi_P + \Phi_e)$

$$\left\{ \begin{array}{l} \nabla_x \sqrt{\rho} + \frac{1}{2} \nabla_x \Phi_P \sqrt{\rho} \in L^2_{t,x}, \\ -\Delta_x \Phi_P = \rho \in L^\infty_t (L^1_x) \\ \nabla_x \Phi_P \in L^\infty_t (L^2_x) \\ \rho = \rho_b \text{ in } H^{1/2}(\partial\omega) \end{array} \right.$$

$$\implies \rho, \nabla_x \sqrt{\rho} \text{ and } \nabla \Phi_P \sqrt{\rho} \in L^2_{t,x}$$

and then

$$J(\rho, \Phi_P + \Phi_e) = -\mathbb{D}(x) [\nabla_x \rho + \rho \nabla_x (\Phi_P + \Phi_e)] \in L^1$$

- Decay of the solution of the cell problem ???

$$\nu \cdot \nabla_y \chi - \nabla_y \Phi_H \cdot \nabla_\nu \chi - Q(\chi) = g$$

- Is it possible to prove a double-scale averaging Lemma ??
- Extension of this analysis to the Fermi-Dirac model ?
- General electric field ? : No explicit relative Maxwellian !
- If σ is not symmetric ?? "No explicit entropy-dissipation"

THANK YOU